Radiative Transfer in Highly Thick Media through Rayleigh and Raman Scattering with Atomic Hydrogen



Seok-Jun Chang

Supervisor : Hee-Won Lee



1.Introduction

2. Radiative Transfer for Raman Scattering

3. Lyα Radiative Transfer for Lyα Blob

4. Future Work

Radiative Transfer

Radiative transfer is the physical phenomenon of energy transfer in the form of electromagnetic radiation.



Scattering with Atomic Hydrogen



Type of Scattering with Atomic Hydrogen



- **Resonance scattering** is the transition between the states.
- **Rayleigh and Raman scattering** occur as the bounded electron excited to the energy level between the states.
- The excited hydrogen atom may de-excite into an excited state resulting in reemission of a lower energy photon, which is called *Raman Scattering.* If the de-excitation is made into the ground state, then result is an elastic scattering, which is also called *Rayleigh scattering.*

Scattering with Atomic Hydrogen



Type of Scattering with Atomic Hydrogen

Literature	Subject	Type of Scattering	Object
Chang et al. 2015	Scattering Wing of $H\alpha$ and $H\beta$	Raman Scattering	AGN
Lee et al. 2016	Raman Scattered O VI	Raman Scattering	Symbiotic Stars
Chang et al. 2017	Polarization of Lyα	Rayleigh scattering	AGN
Chang et al. 2018A	Resonantly Scattered Hα and Lyβ	Resonance Scattering	ISM
Chang et al. 2018B	Scattering Wing of $H\alpha$ and $H\beta$	Raman Scattering Thomson Scattering	Symbiotic Stars
Choi et al. 2020	Raman scattered He II	Raman scattering	Planetary Nebulae
Chang et al. 2020	Grid-Based Simulation	Rayleigh scattering Raman scattering	STaRS
Chang et al. prep	Scattered Lyα in CGM and IGM	Resonance Scattering Rayleigh Scattering	Lyα Blobs

Radiative Transfer for Raman Scattering

Raman Scattering in Astrophysics – Schmid 1989



Fig.1. Raman scattered emission bands in the symbiotic star V1016 Cyg. The spectrum was obtained on the 1.93m telescope at the Observatoire de Haute Provence.



Symbiotic Stars The binary stars composed of hot white dwarf and cool red giant

- Raman scattered lines are relatively broader than other emission lines.
- Schmid (1989) identify Raman Scattered O VI at 6825 Å and 7082 Å in V1016 Cyg.
- The incident photons are O VI λ 1032 and λ 1038.

Raman Scattering in Astrophysics – Schmid 1989



Cyg. The spectrum was obtained on the 1.93m telescope at the Observatoire de Haute Provence.



Fig.2. Schematic energy level diagram for Raman scattering of OVI photons by neutral hydrogen.

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- Schmid (1989) identify Raman Scattered O VI at 6825 Å and 7082 Å in V1016 Cyg.
- The incident photons are O VI λ 1032 and λ 1038.

Raman Scattering in Previous Works

- Chang et al. 2015, Formation of Raman wings near Hα, Hβ, and Paα in AGN unification model
- Chang et al. 2018, Broad Balmer wing in Symbiotic Stars
- Chang & Lee 2020, 3D grid-based Monte Carlo code for Raman scattering

Formation of Raman wings in AGN Unification Model



- In AGN unification model, the type of AGN is determine by the line of sight of one unification model.
- The torus obscure the super massive black hole and broad line region.
- The hydrogen column density of the torus is ~ $10^{23} cm^{-2}$ measured by the hardness of X-ray.

Formation of Raman wings in AGN Unification Model



- The continuum near Ly β and Ly γ convert to Raman wings near H α and H β .
- The widths of Raman wings are proportional to $N_{HI}^{0.5}$.
- Raman wings show the asymmetry, when $N_{HI} \ge 10^{23} cm^{-2}$.

Broad Balmer Wings in Symbiotic Stars



- H α in symbiotic stars shows the broad wing over ~ 1000 km/s.
- Thomson scattering wings show the same width near H α and H β .
- Raman scattering wing near H α is 3 times broader than H β .

Broad Balmer Wings in Symbiotic Stars



H α and H β of S-type Symbiotic Stars, Z And (Chang et al. 2018)

- We got the spectrum of symbiotic stars, Z And, by CFHT.
- We compare with the observation data and simulated wings.
- We conclude that Raman scattering largely contribute to the formation of broad wing in Z And, because Hα wing is broader than Hβ.

STaRS – Poster – PSA-05

Welcome to Seok-Jun's Home

Home About CV & Publication STaRS



STaRS is the code for Radiative Transfer through Raman and Rayleigh Scattering with atomic hydrogen. This code is 3D grid based Monte Carlo simulation tracing each generating photon packet. The information of the photon packet include wavelength, position, and polarization

The paper for STaRS is accepted for publication in JKAS. The link for the manuscript in astro-ph https://arxiv.org/abs/2012.03424

seokjun.weebly.com

The link of Github for STaRS https://github.com/csj607/STaRS

Sejong Radiative Transfer through **Rayleigh and Raman Scattering**





STaRS Gen 2: Sejong Radiative Transfer through Raman and Rayleigh Scattering in Dusty Medium 杉

Seok-Jun Chang¹, Hee-Won Lee¹ and Kwang-II Seon² ¹Department of Physics and Astronomy, Sejong University, Seoul, Korea ²Korea Astronomy and Space Science Institute, Daejeon, Korea

STaRS's Homepac

ission features formed through Raman scattering with atomic hydrogen provide unique and crucial information to probe the distribution and kinematics of a thick neutra egion illuminated by a strong far-ultraviolet radiation source. We introduce a new 3-dimensional Monte-Carlo code to describe the radiative transfer of line photons subjec o Raman and Rayleigh scattering with atomic hydrogen. In our Sejong Radiative Transfer through Raman and Rayleigh Scattering (STaRS) code, the position, direction, wavelength, and polarization of each photon is traced until escape. The thick neutral scattering region is divided into multiple cells. Each cell is characterized by elocity and density, which ensures flexibility of the code in analyzing Raman-scattered features formed in a neutral region with complicated kinematics and density stribution. We are continuously developing STaRS to adopt the absorption and scattering effect by dust. This poster introduces STaRS and its current state and study

I. Abstract



Ly α Radiative Transfer for Ly α Blobs

Lyα Blob



Subaru Image of SSA22-LAB1 (Matsuda et al. 2004)



- High redshift z > 2
- Spatially Extended > 100 kpc
- Luminous Lya 10^{43-44} erg/s
- Over-Dense Region
- Proto Galaxy Cluster
- Circumgalactic Medium in LAB \rightarrow

Intercluster Medium in Galaxy cluster

Lyα Blob



Subaru Image of SSA22-LAB1 (Matsuda et al. 2004)





Lyα spectra by Gas kinematics (Yang et al. 2014B)

 $Ly \alpha \ Radiative \ Transfer$

Lyα Emission Mechanism





Superwind (Galactic wind in M82)



Cold accretion (Goerdt+10)



Radio galaxy, AGN (Reuland+03)

- Shock from Super Wind
- Cold Accretion by Massive Dark Matter Halo
- Photoionization by the source
- Star Formation
- Scattering with Atomic Hydrogen

Polarization of Lyα Blob



- Lyα emission photons are unpolarized.
- Scattering with Atomic hydrogen \rightarrow linear polarization
- Polarized Lyα is observed in Lyα blobs (Prescott et al. 2011, Hayes et al. 2011, Beck et al. 2016, You et al. 2017, Kim et al. 2020)

 $Ly\alpha$ Radiative Transfer

Polarization by Scattering





The degree of polarization of the photon after Rayleigh scattering maintain (forward and backward scattering) or increase (90^o scattering).

Lyα Radiative Transfer

1990

Neufeld 1990, 1991 – Analytic Solution

Ahn & Lee 1998, 2002 – First Numerical Method for polarization of scattered Lyα, No spatial distribution

2000

Zheng & Miralda-Escude 2002 – Spectrum & surface brightness

Verhamm et al. 2006 – Spectrum in dusty medium

Dijkstra & Loeb 2008 – Spatial distribution of Polarization, Only Rayleigh scattering

2010 - current

Gronke et al. 2016, 2017 – Spectrum in clumpy medium

Chang et al. 2017 – Polarized Lyα in AGN unification model

Eide et al. 2018 – Polarization of Lyα for LAE, Small physical scale

Seon & Kim 2020 – WF effect with hyper-find structure

Seon et al. in prep – Polarization adopting Stokes parameter, Narrow emission line width

Radiative Transfer for LAB



Can the Lyα point source become LAB through only scattering by an atomic hydrogen in H I halo ?

Essential Components for LAB

- Polarization behavior by core and wing scattering
- Large scale halo ~ 100 kpc
- Broad Lyα emission
- Kinematics of scattering medium

Scattered-Only Lya in Spherical H I halo



$$v(r) = v_{exp} \frac{r}{R_H}$$
 $N_{HI} = \int_{0}^{R_H} \rho_{HI}(r) dr \quad \rho_{HI}(r) = n_0 e^{r/R_e}$

Parameter	Range	Note
$N_{ m HI}$	$10^{18-21} \mathrm{cm}^{-2}$	H I column density
R_e	$0.1R_H - R_H, \infty$	Effective radius
$v_{ m exp}$	$0 - 400 \mathrm{km s^{-1}}$	Expanding velocity
$\sigma_{ m src}$	$100 - 400 \mathrm{km s^{-1}}$	Ly α source width

$$z = 3 (7.855 \text{ kpc/"})$$
 $L_{Ly\alpha} = 10^{44} erg/s$

- We develop the simulation to study scattered-only Lyα.
- The simulation is based on LaRT (Seon & Kim 2020, Seon et al. prep).
- We consider H I spherical halo with $R_H = 100 \text{ kpc}$ surrounding Ly α point source.
- The width of Ly α emission from the source is 100 400 km/s.



- The surface brightness becomes more extended with increasing N_{HI} .
- The degree of polarization is not monotonic.

Lyα Projected Image

Dependence on N_{HI}

Surface brightness Degree of polarization Spectrum 10⁻¹⁵ 60 0.5 Outflow - SFG logN_{HI} [cm Outflow - SFG Outflow - SFG logN_{HI} [cm] logN_{HI} [cm SB [erg s⁻¹ cm⁻² arcsec⁻²] 50 $= 100 \ km \ s^{-1}$ 18 0.4 10⁻¹⁶ 19 20 Relative Flux 700 Relative Flux 40 DoP [%] 30 10^{-17} 20 10⁻¹⁸ 10 0.1 100 -1000-500500 1000 1500 2000 20 40 60 80 40 60 80 100 0 20 $\Delta V [km s^{-1}]$ R_p [kpc] R_p [kpc] 10⁻¹⁵ 0.6 60 Outflow - AGN Outflow - AGN logN_{HI} [cm logN_{HI} [cm logN_{HI} [cm⁻² Outflow – AGN SB [erg s⁻¹ cm⁻² arcsec⁻²] 0.5 50 $\sigma_{src} = 400 \ km \ s^{-1}$ 18 10⁻¹⁶ 19 6.48.68.68.78.79.7<l 40 20DoP [%] 30 10⁻¹⁷ 20 10⁻¹⁸ 10 0.1 0 $500 \\ \Delta V [km s⁻¹]$ -5001000 20 100 -10001500 2000 0 20 40 60 80 100 0 40 60 80 0 R_p [kpc] R_p [kpc] SB is more extended $N_{HI} \uparrow \rightarrow$ **Complex polarization behavior** Spectrum is more broadened



- At $N_{HI} = 10^{18-19} cm^{-2}$, the photon easily escape through several wing scattering.
- At $N_{HI} = 10^{18} cm^{-2}$, the degree of polarization decrease due to resonance scattering.
- At $N_{HI} = 10^{21} cm^{-2}$, the photon must go through multiple wing scattering.

Dependence on v_{exp} in High N_{HI} (10²¹ cm⁻²) Halo



Dependence on v_{exp} in Low N_{HI} (10¹⁹ cm⁻²) Halo



Size of Lyα Halo





- When the polarization is maximum, the effect of the single wing scattering is strongest.
- At $N_{HI} = 10^{18} cm^{-2}$, the polarization decrease with increasing v_{exp} .
- At $N_{HI} = 10^{21} cm^{-2}$, the polarization behavior does not depend on σ_{src} .

Spectral Peak at ΔV_{peak}



- ΔV_{peak} does not depend on v_{exp} in $N_{HI} < 10^{20} \ cm^{-2}$.
- ΔV_{peak} does not depend on σ_{src} .
- The multiple wing scattering causes the variation of ΔV_{peak} as v_{exp} increases.



- In this work, the simulation for $Ly\alpha$ radiative transfer is based on LaRT.
- We test that scattered-only Lyα becomes LAB in spherical H I halo.
- The surface brightness profiles become more extended with increasing N_{HI} and v_{exp} .
- The polarization pattern is always concentric because of spherical scattering region.
- The degree of polarization is not monotonic due to resonance scattering in low N_{HI} regime.
- When $N_{HI} \ge 10^{20} cm^{-2}$, Ly α halo is extended over 100 kpc.

Future Work

Lyα Radiative Transfer in Clumpy Medium

Model C (Clumpy Medium)



Clumpy Medium

- Gronke et al. 2017 studied the Lya spectrum profile in clumpy medium.
- They find that Lya spectrum in a clumpy medium becomes similar to the spectra of

a continuous medium.

We revisits the results in Gronke et al.
2017 using our simulation.



Thanks